

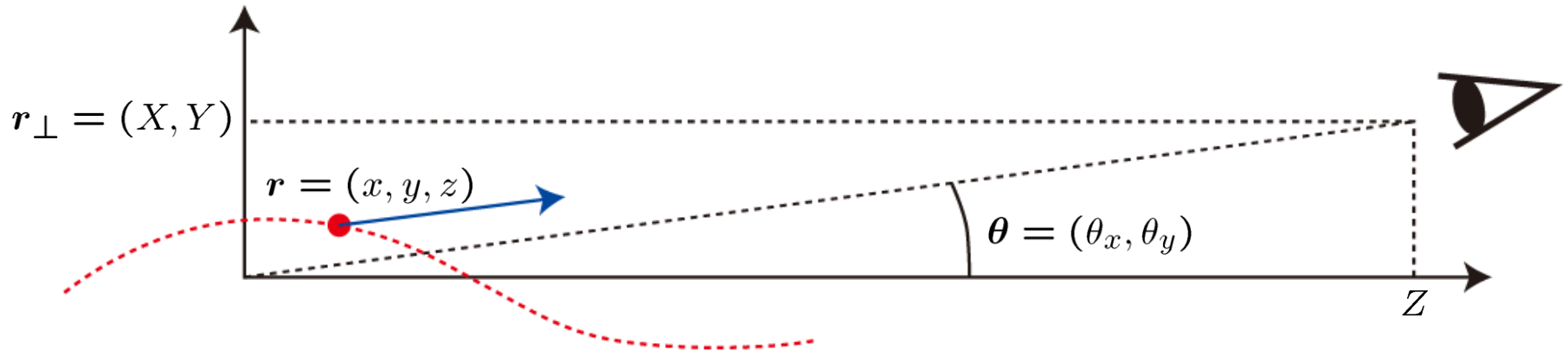
# Light Source II

*Takashi TANAKA (RIKEN SPring-8 Center)*

## Characteristics of SR (2)

- Electron Trajectory in the ID
- Qualitative Description of Wiggler Radiation
- Qualitative Description of Undulator Radiation

# Coordinate Systems



SR emitted by an electron moving at  $\mathbf{r} = (x, y, z)$   
 Observation of SR at  $\mathbf{R} = (X, Y, Z)$

If the far-field approximation ( $|\mathbf{r}| \ll Z$ ) is applicable, the radiation pattern depends only on the observation angle  $\theta = (\theta_x, \theta_y)$ .

# Field Integrals

$$\frac{d\mathbf{P}}{dt} = m\gamma \frac{d\mathbf{v}}{dt} = -e\mathbf{v} \times \mathbf{B} \rightarrow \begin{cases} m\gamma \dot{v}_x = -e(v_y B_z - v_z B_y) \\ m\gamma \dot{v}_y = -e(v_z B_x + v_x B_z) \end{cases}$$

Equation of motion of an electron moving in a magnetic field  $\mathbf{B}$

$$\downarrow B_z \equiv 0$$

$$m\gamma \frac{dv_{x,y}}{v_z dt} = m\gamma \frac{dv_{x,y}}{dz} = \pm e B_{y,x}$$

$$\begin{aligned} \beta_{x,y} &= \pm \frac{e}{\gamma m c} \int^z B_{y,x}(z') dz' \equiv \pm \frac{e}{\gamma m c} I_{1y,1x}(z) \\ x, y &= \pm \frac{e}{\gamma m c} \int \int^{z'} B_{y,x}(z'') dz'' \equiv \pm \frac{e}{\gamma m c} I_{2y,2x}(z) \end{aligned}$$

$I_1, I_2$  : 1st and 2nd field integrals of the ID

# Trajectory in an Ideal ID

$$\left\{ \begin{array}{l} B_x(z) = 0 \\ B_y(z) \sim B_0 \sin\left(\frac{2\pi z}{\lambda_u}\right) \end{array} \right\} \quad \left\{ \begin{array}{l} \beta_y = 0 \\ \beta_x = \frac{K}{\gamma} \cos\left(\frac{2\pi z}{\lambda_u}\right) \end{array} \right\} \quad \left\{ \begin{array}{l} y = 0 \\ x = \frac{\lambda_u K}{2\pi\gamma} \sin\left(\frac{2\pi z}{\lambda_u}\right) \end{array} \right\}$$

magnetic field



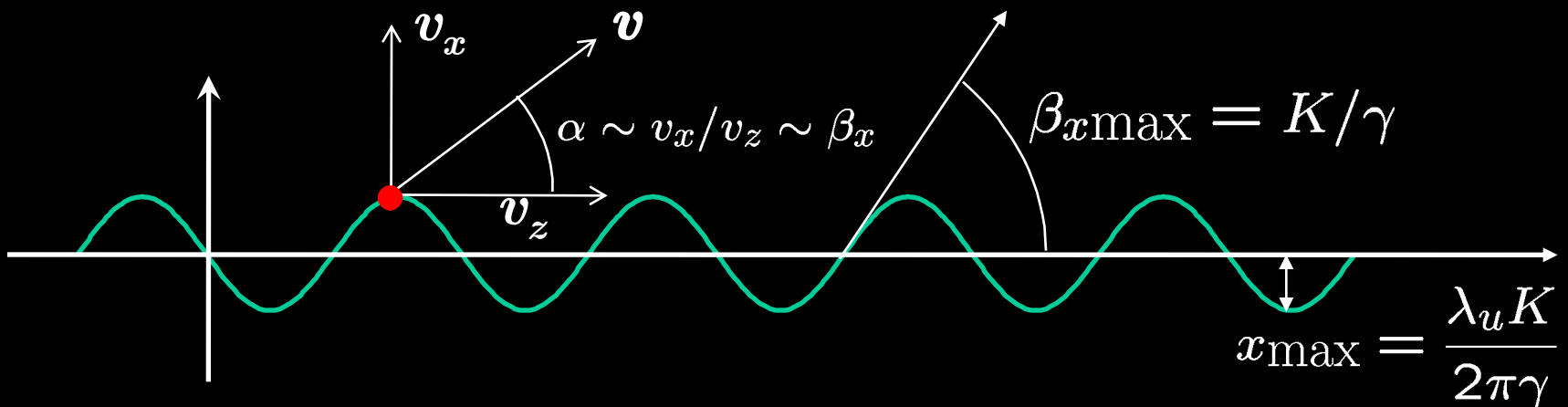
velocity



position

$$K = \frac{eB_0\lambda_u}{2\pi mc} = 93.37 B_0(\text{T})\lambda_u(\text{cm})$$

K value, Deflection parameter



$$E_e = 8\text{GeV}, K=1, \lambda_u = 5\text{cm} : \beta_{x\text{max}} = 64\mu\text{rad}, x_{\text{max}} = 0.5\mu\text{m}$$

# Effects due to the ID Magnetic Field

transverse  
velocity

$$\beta_x = \frac{K}{\gamma} \cos\left(\frac{2\pi z}{\lambda_u}\right)$$



longitudinal  
velocity

$$\beta_z = \sqrt{\beta^2 - \beta_x^2}$$

$$= \underbrace{1 - \frac{1}{2\gamma^2} - \frac{K^2}{4\gamma^2}}_{\bar{\beta}_z: \text{average velocity}} - \underbrace{\frac{K^2}{4\gamma^2} \cos\left(\frac{4\pi z}{\lambda_u}\right)}_{\text{oscillating term}}$$

$\bar{\beta}_z$ : average velocity

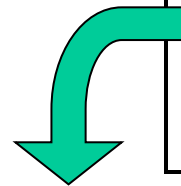
oscillating term

ID field induces:

- transverse(x) oscillation
- longitudinal (z) oscillation
- effective deceleration( $\Delta\beta_z = K^2/4\gamma^2$ )

# General Form of Time Squeezing

$$\frac{d\tau}{dt} = 1 - \boldsymbol{\beta} \cdot \mathbf{n}$$



$$\begin{aligned}\beta_z &= \sqrt{\beta^2 - \beta_x^2 - \beta_y^2} \\ &\sim 1 - (\gamma^{-2} + \beta_x^2 + \beta_y^2)/2 \\ n_z &\sim 1 - (\theta_x^2 + \theta_y^2)/2\end{aligned}$$

$$= \frac{1}{2\gamma^2} + (\theta_x - \beta_x)^2 + (\theta_y - \beta_y)^2$$

Time squeezing takes place most significantly when the direction of the electron motion coincides with that of observation ( $\boldsymbol{\beta} = \boldsymbol{\theta}$ ).

## Characteristics of SR (2)

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- Qualitative Description of Wiggler Radiation
- Qualitative Description of Undulator Radiation



# Wiggler Radiation

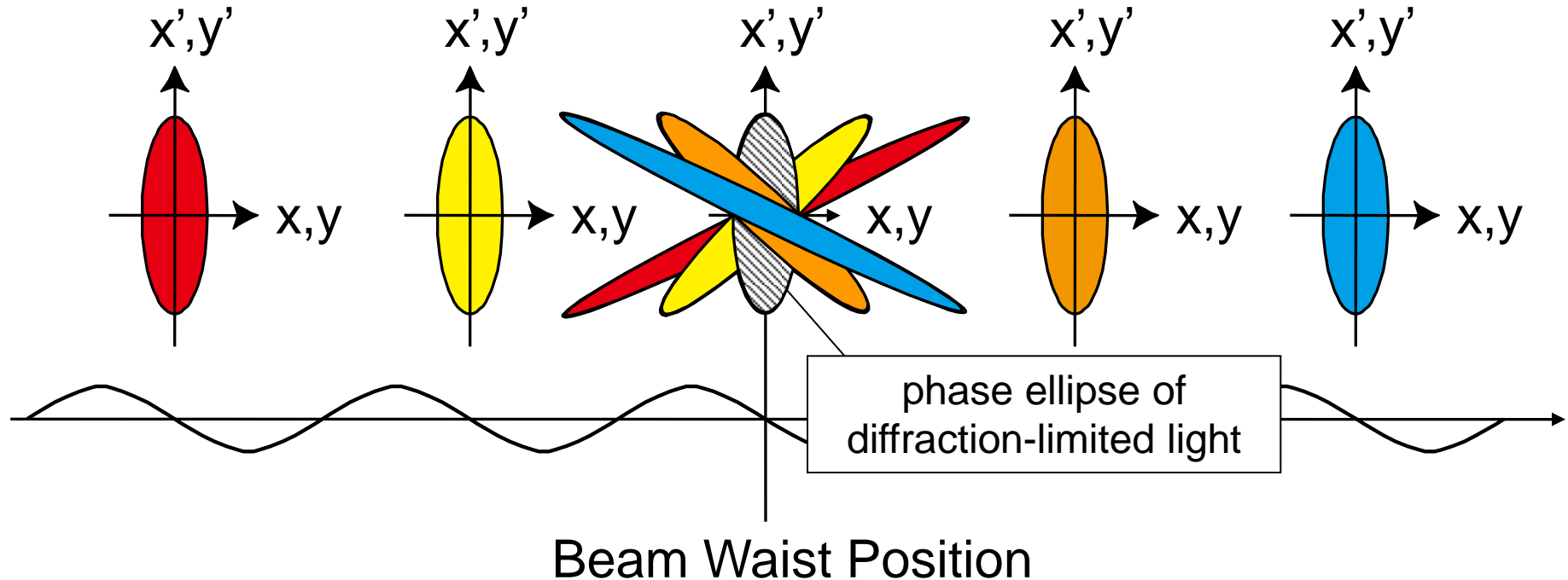
- Wiggler radiation (WR) is regarded as **incoherent sum of SR** at each position.
  - Summation as photons in the framework of geometrical optics.

$$\text{Flux: } F_W \sim 2N F_{BM}$$

$$\text{Emittance: } \sigma_{x',y'} \times \sigma_{x,y} \gg \lambda/4\pi$$

$$\text{Brilliance: } B_W \ll 2N B_{BM}$$

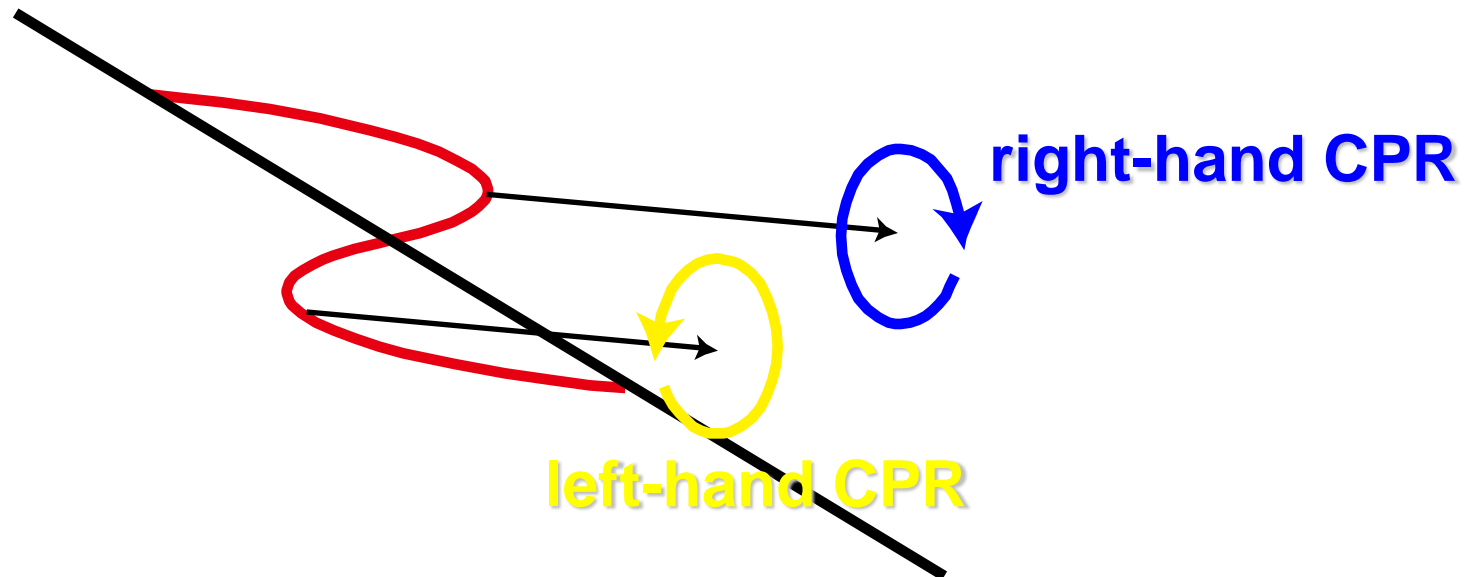
# Photon Distribution in Phase Space



- Larger  $N$  results in larger area of photon distribution in the phase space, i.e., larger emittance.
- $B$  does not linearly depend on  $N$

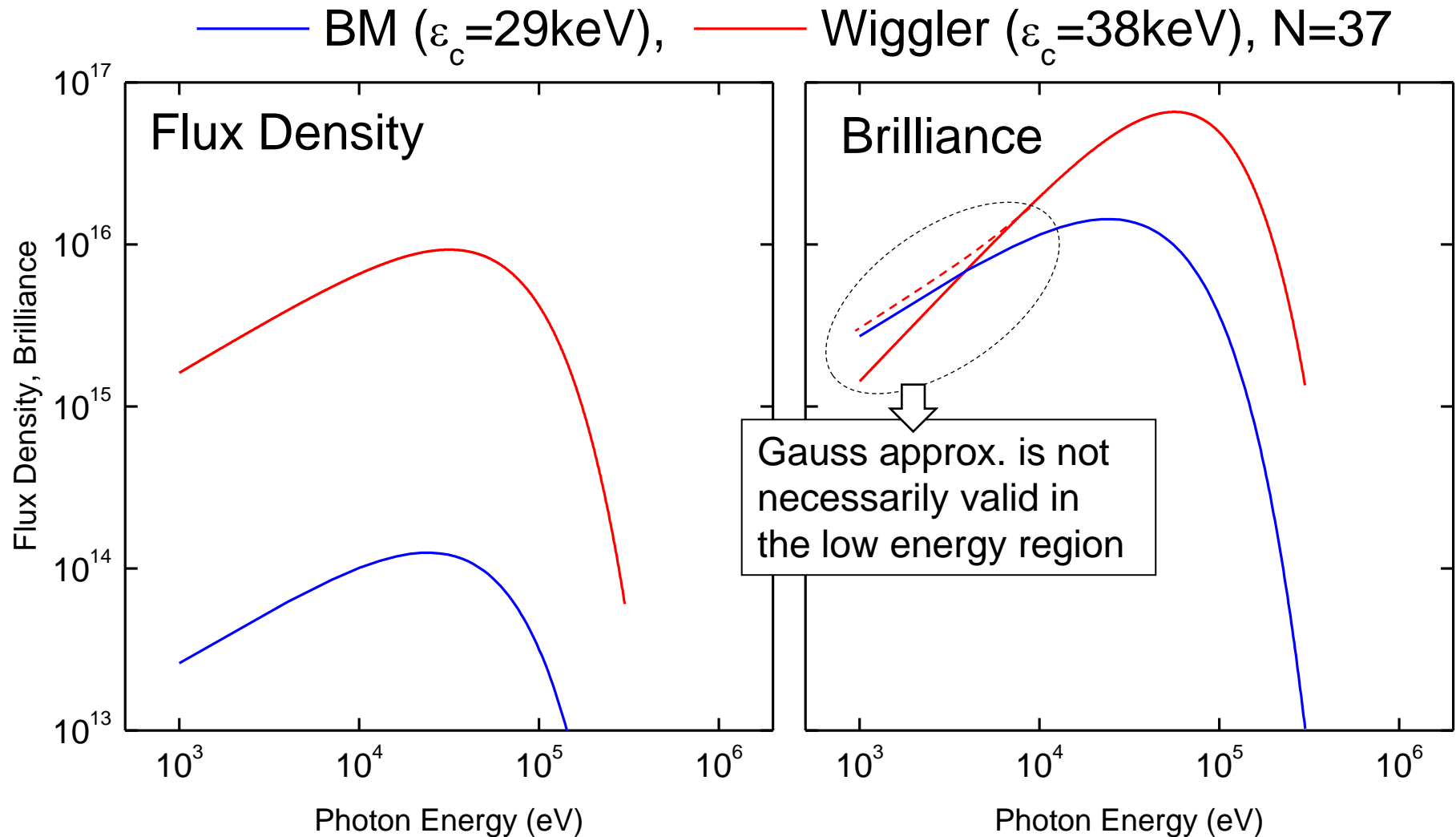
# Polarization

- No circular polarized radiation (CPR) is observed unlike the BM radiation even off axis, due to cancellation of CPR components.



- EMPW is a special wiggler to utilize CPR by introducing a vertical motion.

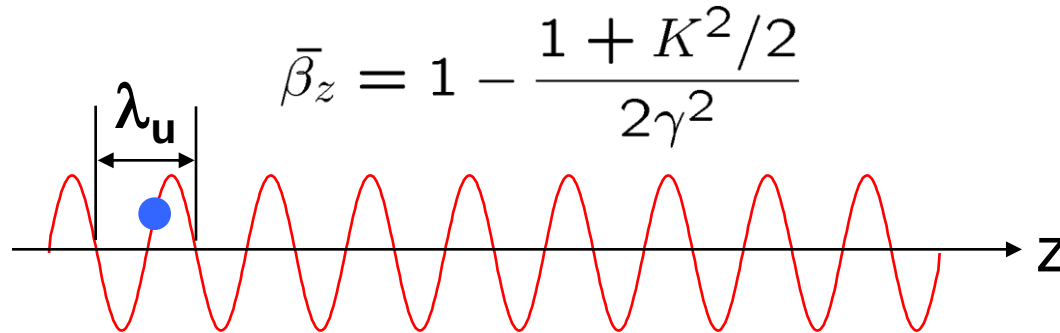
# Comparison with BM Radiation



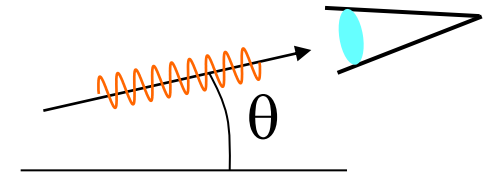
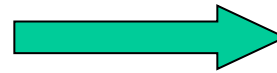
## Characteristics of SR (2)

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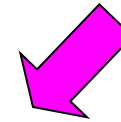
# Fundamental Wavelength



$T = \lambda_u / v_z = \lambda_u / c$   
 period of electron motion  
 = period of emitted light

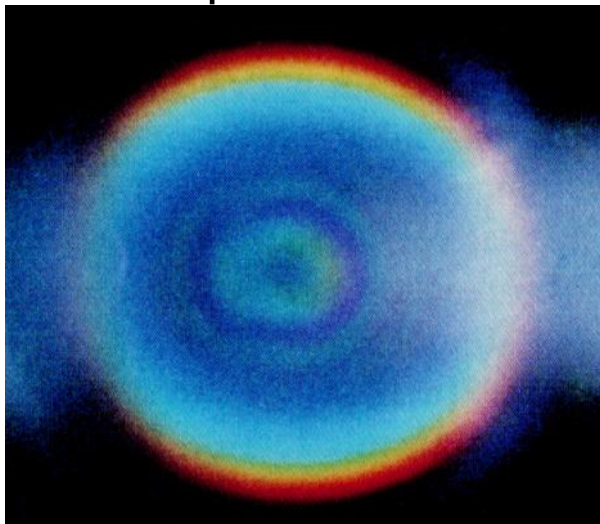


$T' = T(1 - \bar{\beta}_z \cos \theta)$   
 period of observed light



Fundamental Wavelength  $\lambda_1$

$$\begin{aligned}
 \lambda_1 &= \lambda_u (1 - \bar{\beta}_z \cos \theta) \\
 &= \frac{\lambda_u}{2\gamma^2} (1 + \gamma^2 \theta^2 + K^2/2) \\
 \omega_1 &= 2\pi c / \lambda_1
 \end{aligned}$$

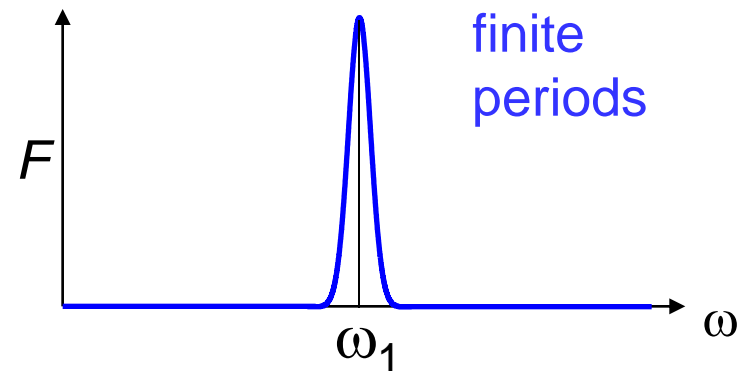
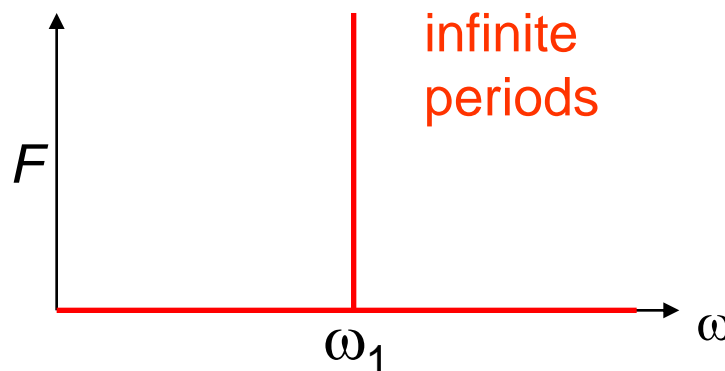


# UR with Infinite Periods

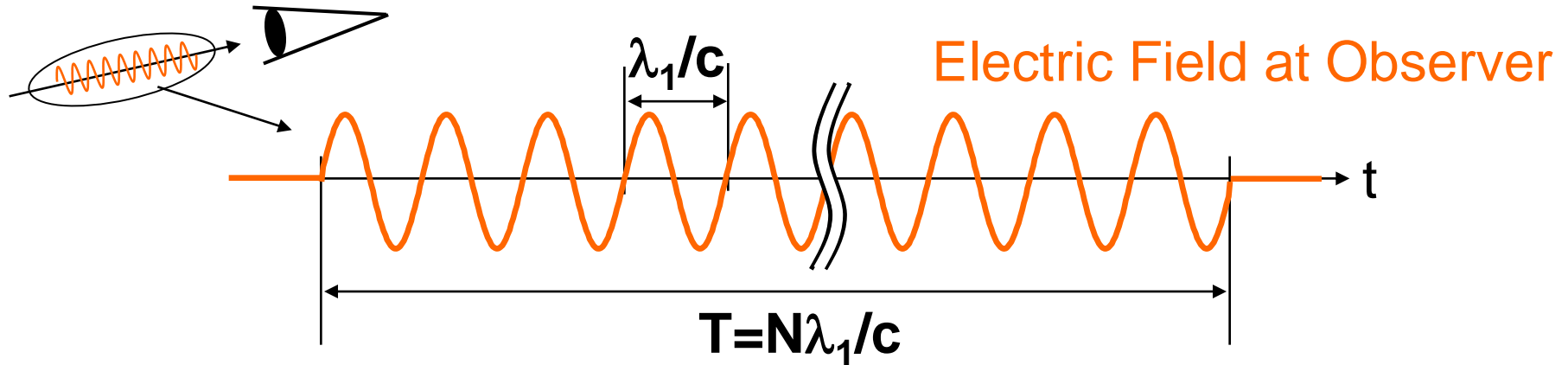
- If the undulator length is infinite, the pulse duration is infinitely long, and thus the radiation is completely monochromatic with line spectrum.

$$\frac{d^2 F}{dx' dy'} \propto \delta(\omega - \omega_1) = \delta\left(\omega - \frac{4\pi c \gamma^2 / \lambda_u}{1 + K^2/2 + \gamma^2 \theta^2}\right)$$

- In practice, the undulator length is finite, so the line spectrum is broadened.



# Effects due to Finite Periods



$$E(t) = \begin{cases} E_0 \sin \omega_1 t & ; -T/2 \leq t \leq T/2 \\ 0 & ; t < -T/2, T/2 < t \end{cases}, \quad \omega_1 = 2\pi c/\lambda_1$$

Fourier Transform

$$\frac{d^2 F}{dx' dy'} \propto |\tilde{E}(\omega)|^2 \propto \text{sinc}^2 \left[ \pi N \frac{\omega - \omega_1(\theta)}{\omega_1(\theta)} \right]$$

Square of “sinc” function dominates the UR



# Brief Note on UR Formulae

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- In the previous derivations of UR spectral function, no knowledge on electrodynamics is required.
- In practice,  $E_0$  is a complicated function of  $\theta$  and  $K$ , and needs to be calculated by Fourier transforming the electric field derived from the Lienard-Wiechert potential.
- However, the simple derivation gives us a clear understanding on UR properties.

# Energy and Angular Profile of UR

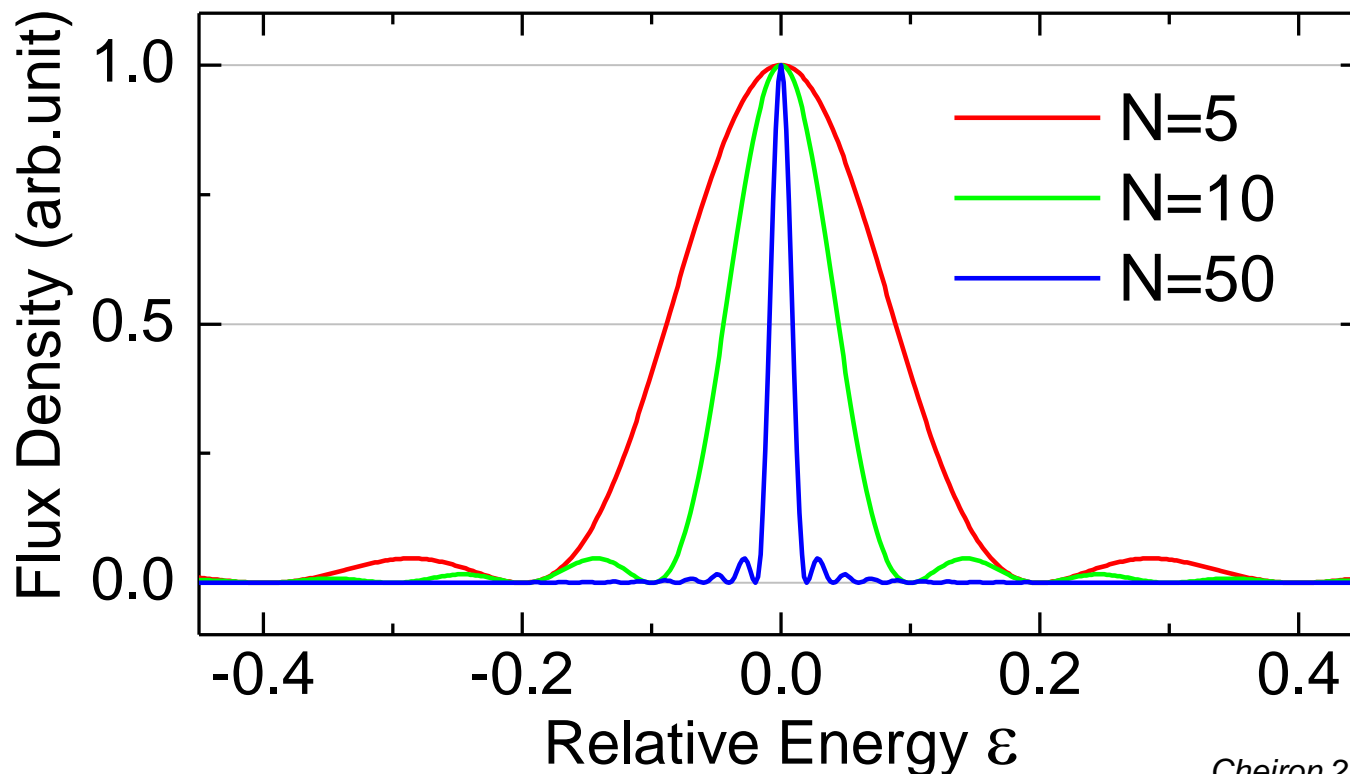
$$\frac{d^2 F(\omega, \theta)}{d\Omega d\omega / \omega} = F_0 \text{sinc}^2 \left[ \pi N \frac{\omega - \omega_1(\theta)}{\omega_1(\theta)} \right]$$

$$= \begin{cases} \text{Energy Profile at } \theta = 0 \\ F_0 \text{sinc}^2(N\pi\varepsilon) \\ \quad ; \varepsilon = [\omega - \omega_1(0)]/\omega_1(0) \\ \\ \text{Angular Profile at } \omega = \alpha\omega_1(0) \\ F_0 \text{sinc}^2[N\pi(\alpha\Theta^2 + \alpha - 1)] \\ \quad ; \Theta = \gamma\theta/\sqrt{1 + K^2/2} \end{cases}$$

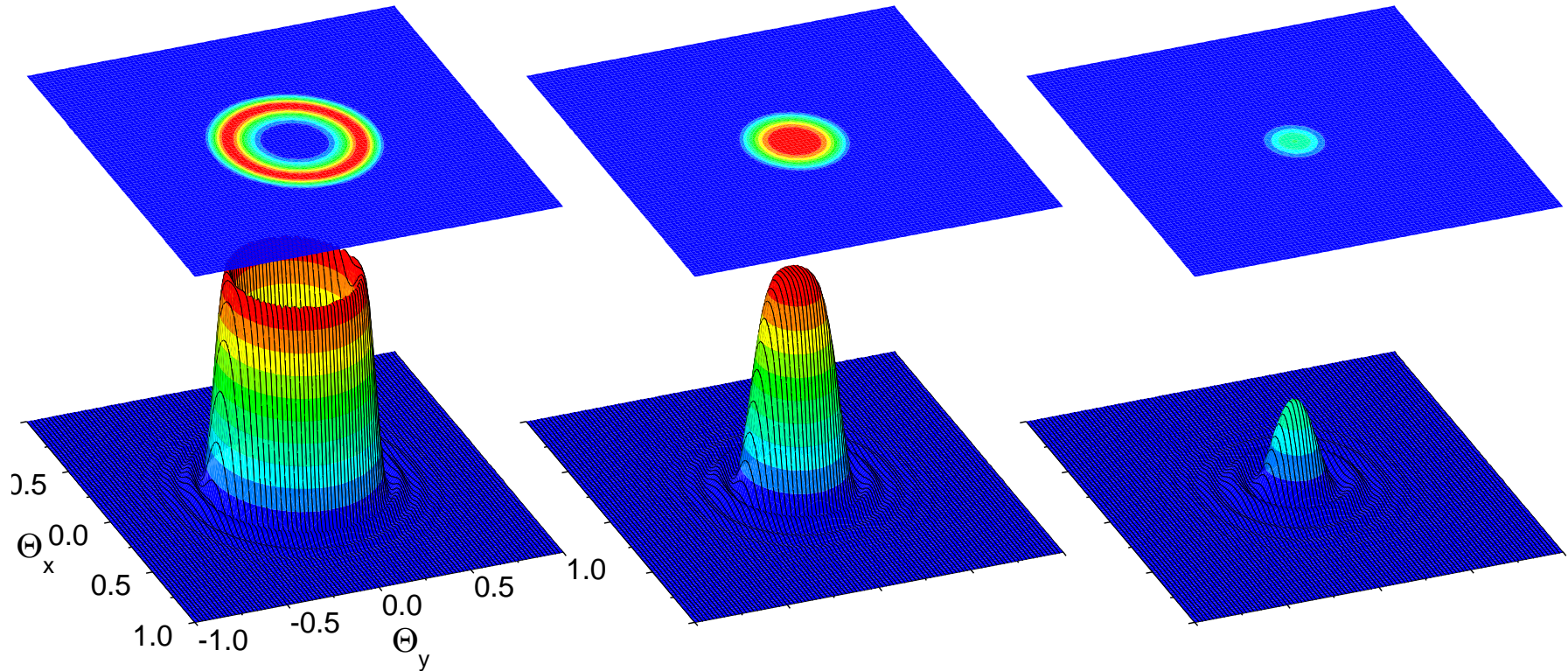
# Energy Profile: Example

$$\frac{d^2 F}{dx' dy'} = F_0 \text{sinc}^2(N\pi\varepsilon); \quad \text{sinc}^2(2.783) \sim 1/2$$

$$\xrightarrow{\quad} \left. \frac{\Delta\omega}{\omega_1(0)} \right|_{FWHM} \sim \frac{0.8858}{N}$$



# Angular Profile: Example



$$\omega = 0.9\omega_1(0)$$

lower energy

$$\omega = \omega_1(0)$$

fundamental  
energy

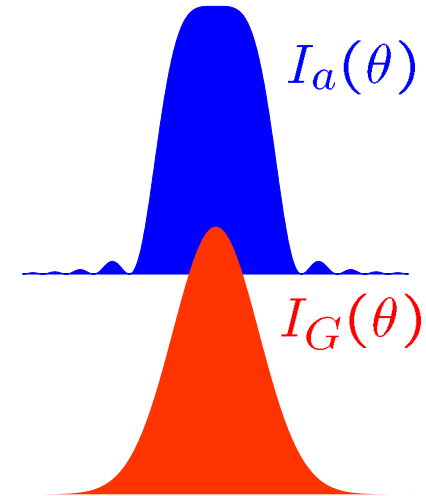
$$\omega = 1.05\omega_1(0)$$

higher energy

# Angular Divergence and Beam Size

Angular Profile at  $\omega=\omega_1(0)$

$$I_a(\theta) = F_0 \text{sinc}^2 \left[ \frac{\pi N (\gamma \theta)^2}{1 + K^2/2} \right]$$



Gaussian Profile with  $\sigma_{r'}$   
 $I_G(\theta) = F_0 \exp(-\theta^2 / 2\sigma_{r'}^2)$

approximation

$$\sigma_{r'} = \sqrt{\frac{1 + K^2/2}{4N\gamma^2}} = \sqrt{\frac{\lambda_1}{2L}}$$

Angular Divergence  
of UR ( $L=N\lambda_u$ )

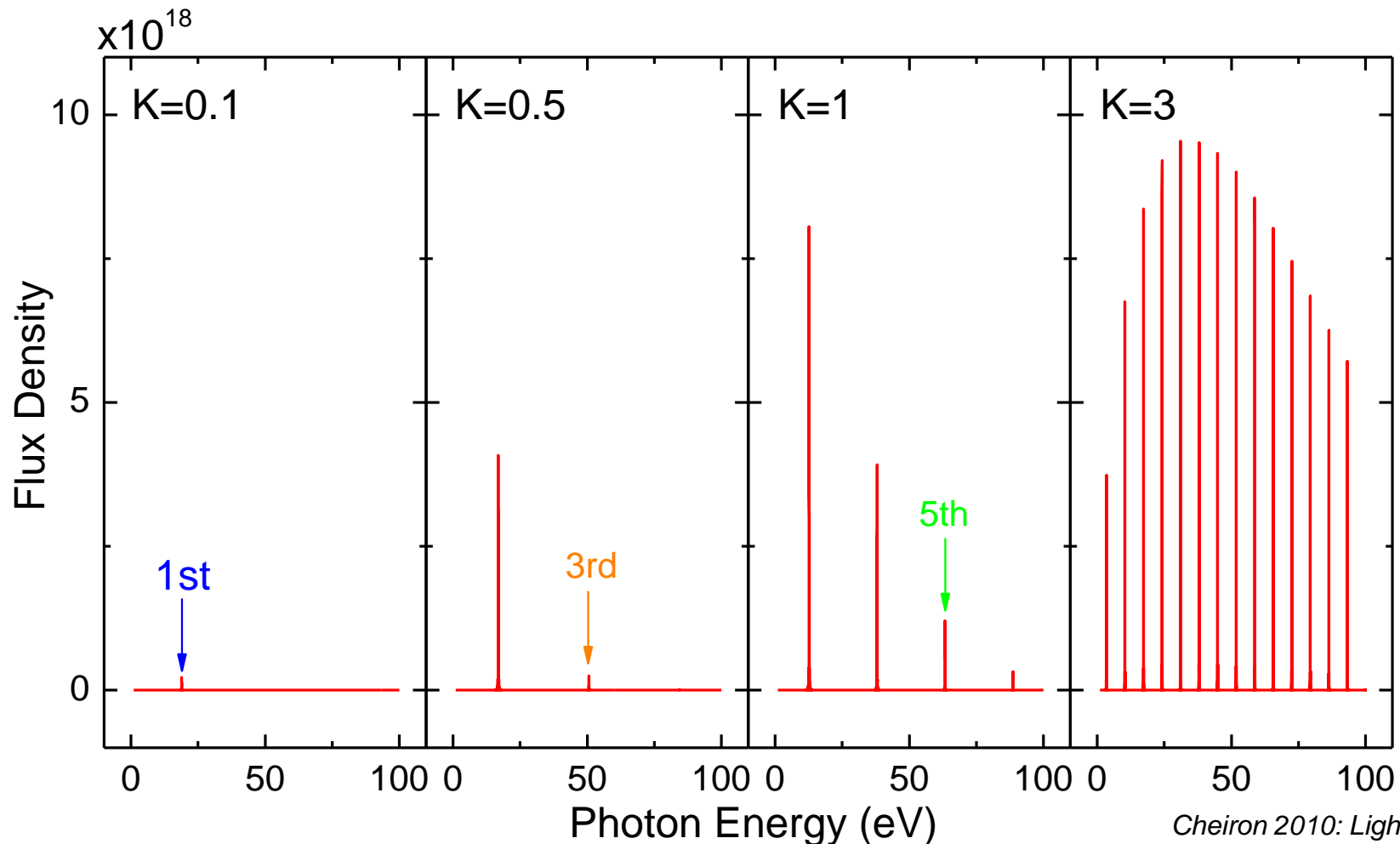
Diffraction Limit (UR is  
Spatially Coherent)

$$\sigma_r = \frac{\lambda_1}{4\pi\sigma_{r'}} = \frac{\sqrt{\lambda_1 L}}{4\pi}$$

Beam Size of UR

# Higher Harmonics

- In addition to  $\omega_1$ , photons with the energy at  $n\omega_1$  is also observed, where  $n$  is an integer.



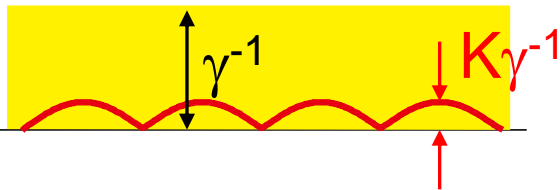
# Mechanisms of Higher Harmonics

 $K \ll 1$ 

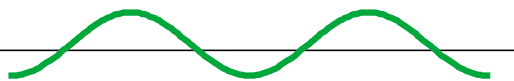
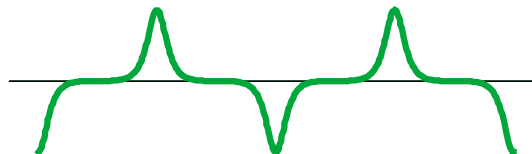
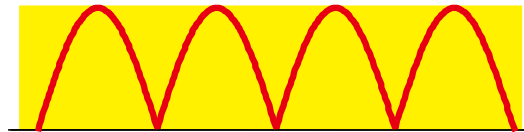
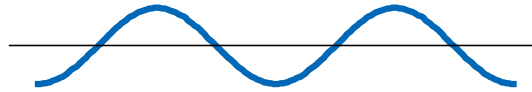
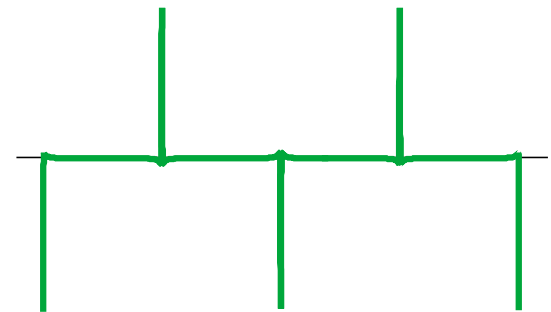
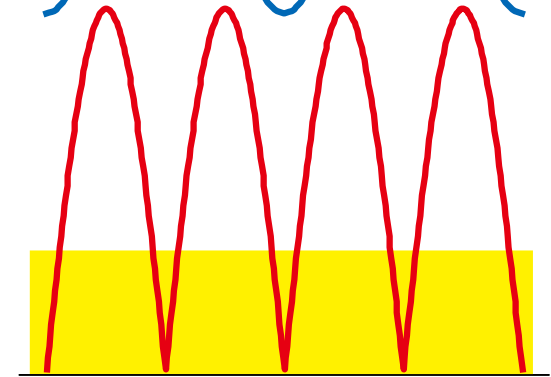
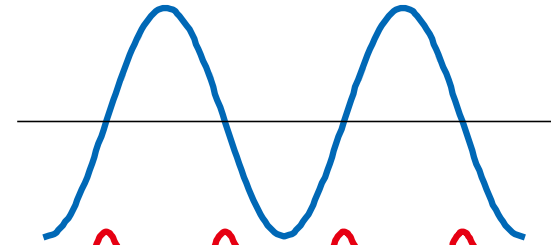
Electron Orbit



Deflection Angle



Radiation E-field


 $K \sim 1$ 

 $K \gg 1$ 


# Optical Properties of Higher Harmonics

For the n-th harmonic radiation,

$$\frac{d^2 F}{dx' dy'} = F_0 \text{sinc}^2 \left[ \pi n N \frac{\omega - n\omega_1(\theta)}{n\omega_1(\theta)} \right]$$



$$\left. \frac{\Delta\omega}{n\omega_1(0)} \right|_{FWHM} \sim \frac{0.8858}{nN} \quad \text{band width}$$

$$\sigma_{r'n} = \sqrt{\frac{1 + K^2/2}{4nN\gamma^2}} = \sqrt{\frac{\lambda_1/n}{2L}} \quad \text{angular divergence}$$

$$\sigma_{rn} = \frac{\lambda_1/n}{4\pi\sigma_{r'n}} = \frac{\sqrt{L\lambda_1/n}}{4\pi} \quad \text{beam size}$$



# Practical Knowledge on Undulator Radiation

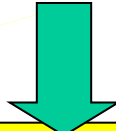
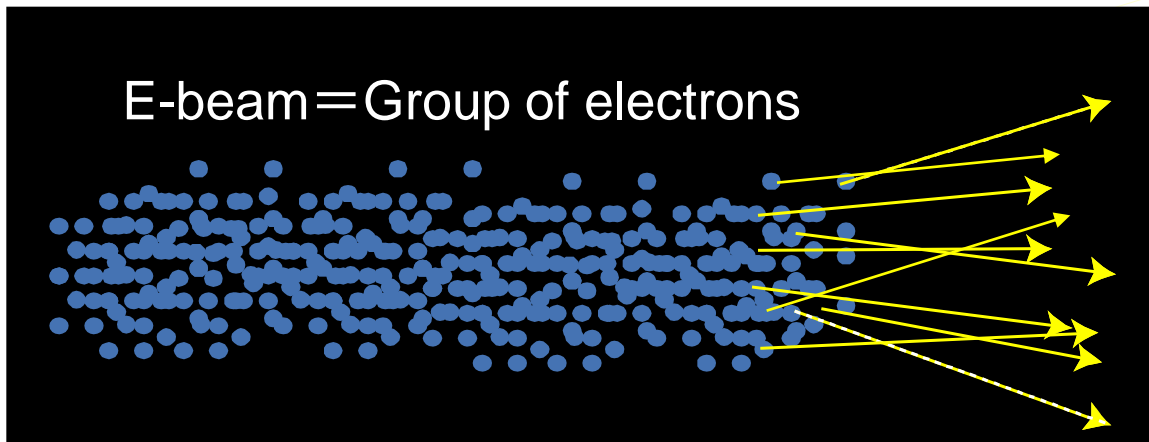
# “Practical Knowledge” on UR

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- The undulator is a “double-edged sword”!
  - Much higher brilliance is available than the BM and wiggler, *if appropriately used*.
  - High heat load on optical elements
  - Quasi-monochromatic and small angular spread imposes an accurate adjustment of BL components.
- Practical knowledge for the utilization of UR
  - Effects due to the electron beam quality
  - Simple evaluation of optical properties
  - Heat load reduction
  - .....

# Electron Beam Quality (1)

- The property of UR from a single electron is similar to a laser
  - Small Size & Angular Divergence, Narrow Bandwidth
- In practice, UR in the beamline is emitted by the beam comprising a huge number of electrons.
- These electrons have different positions, angle, and energies.



The performance of UR depends largely on the e-beam quality

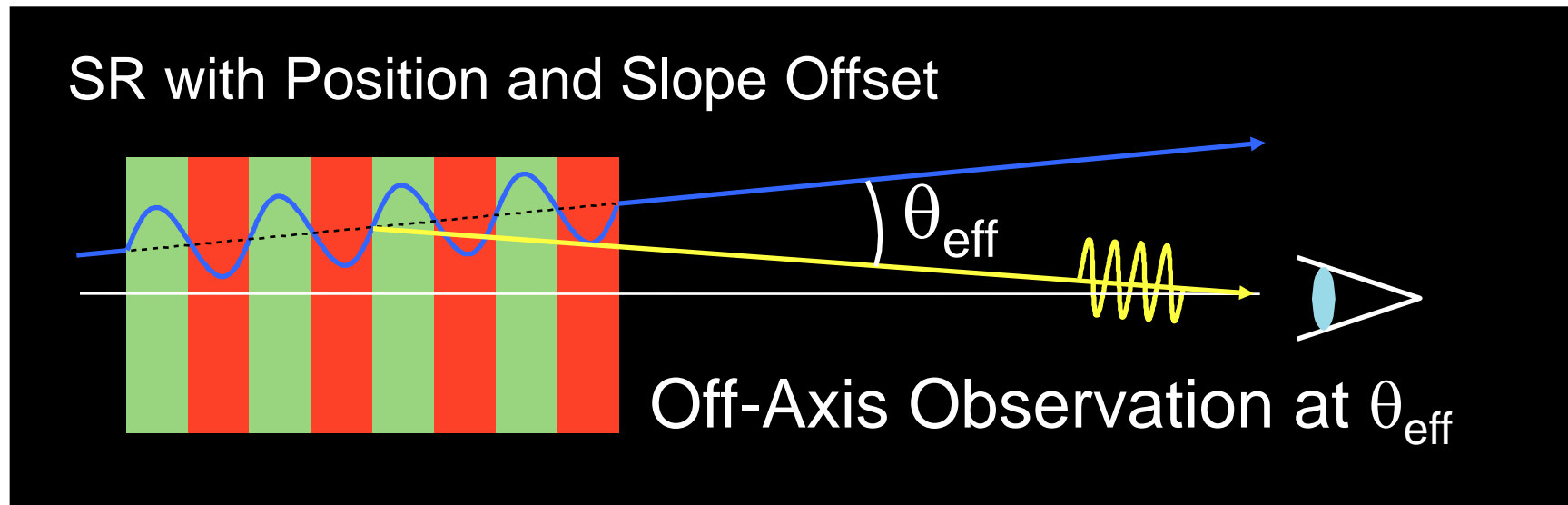
# Electron Beam Quality (2)

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- Electron-beam parameters related to the performance of UR
  - Beam Size  $\sigma_x, \sigma_y \rightarrow$  Source Size
  - Angular Divergence  $\sigma_{x'}, \sigma_{y'} \rightarrow$  Directivity
    - The minimum value of the product  $\sigma_x \sigma_{x'}$  and  $\sigma_y \sigma_{y'}$  are called the electron beam emittances in the x and y directions.
  - Energy Spread  $\sigma_E/E \rightarrow$  Monochromaticity

# Effects due to Finite Emittance (1)

- Effects due to Finite Emittance of the Electron Beam
  - Injection to the undulator with angular and positional offset



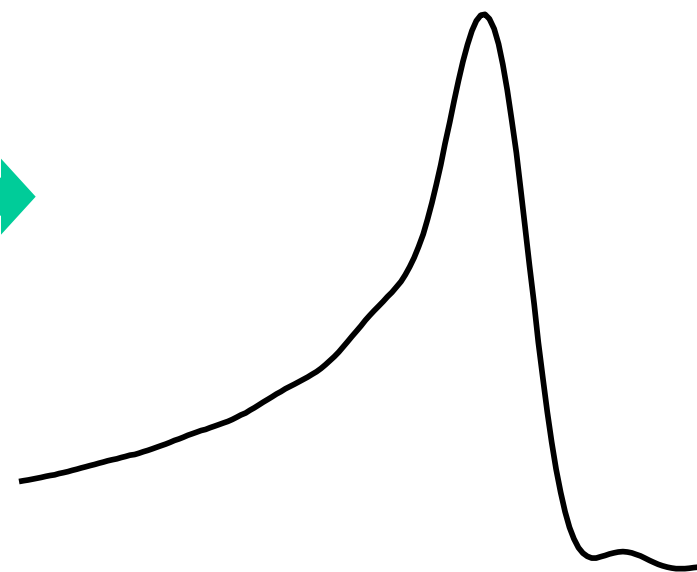
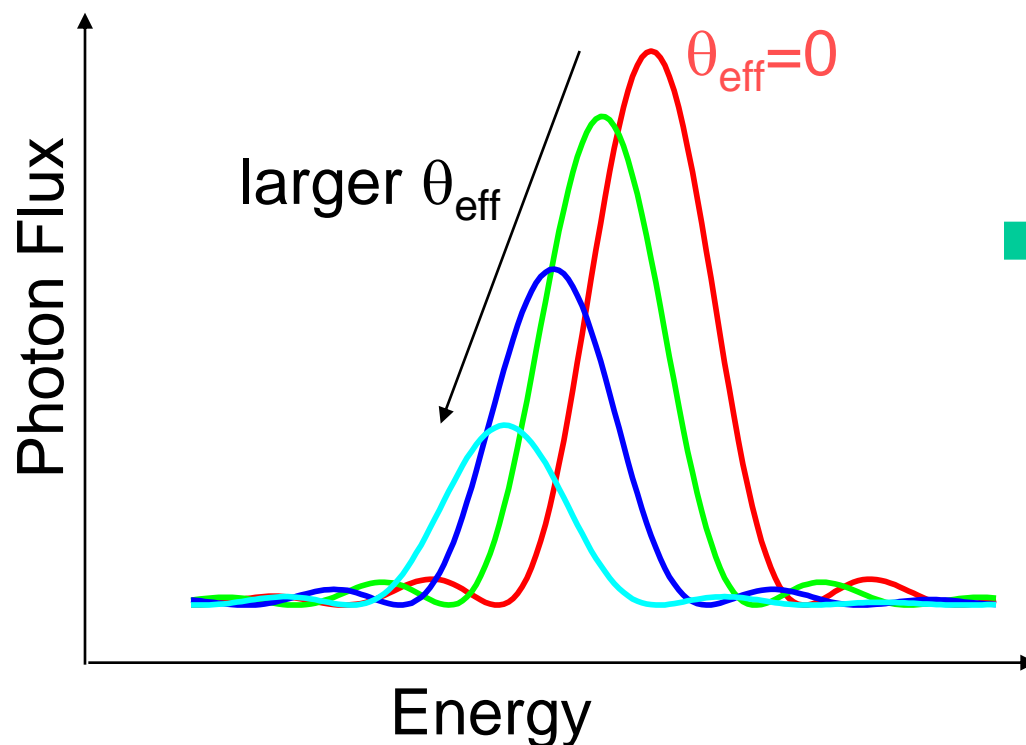
# Effects due to Finite Emittance (2)

Off-axis observation at  $\theta_{\text{eff}}$



Peak shift to  
lower energy

$$\omega_1(\theta) = \frac{4\pi c \gamma^2 / \lambda_u}{1 + \boxed{\gamma^2 \theta^2} + K^2/2}$$



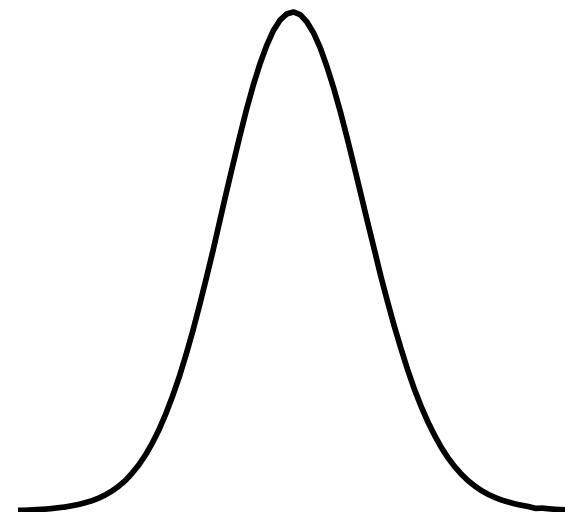
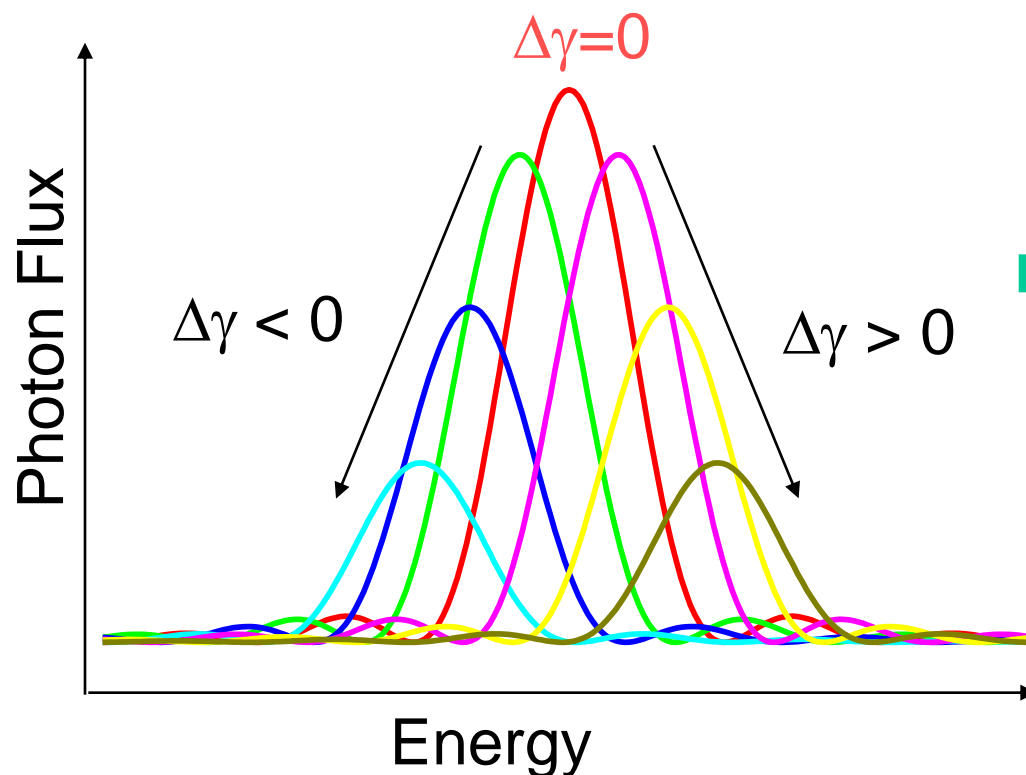
# Effects due to the Energy Spread

Electron with an offset of  $\Delta\gamma$

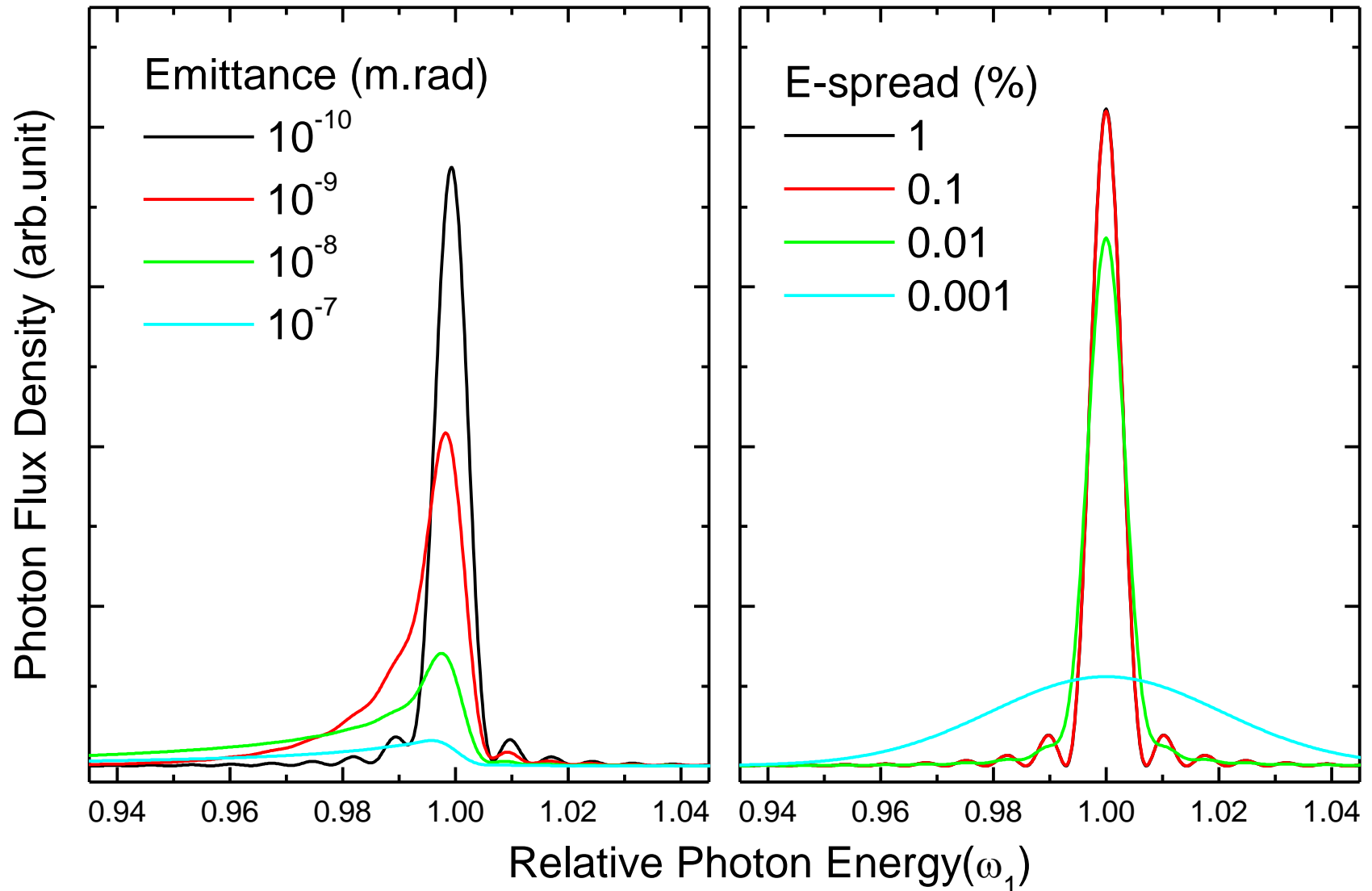


Energy shift of  $\omega_1$

$$\omega_1(\gamma) = \frac{4\pi c \gamma^2 / \lambda_u}{1 + K^2/2}$$



# Examples

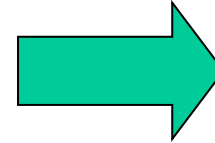




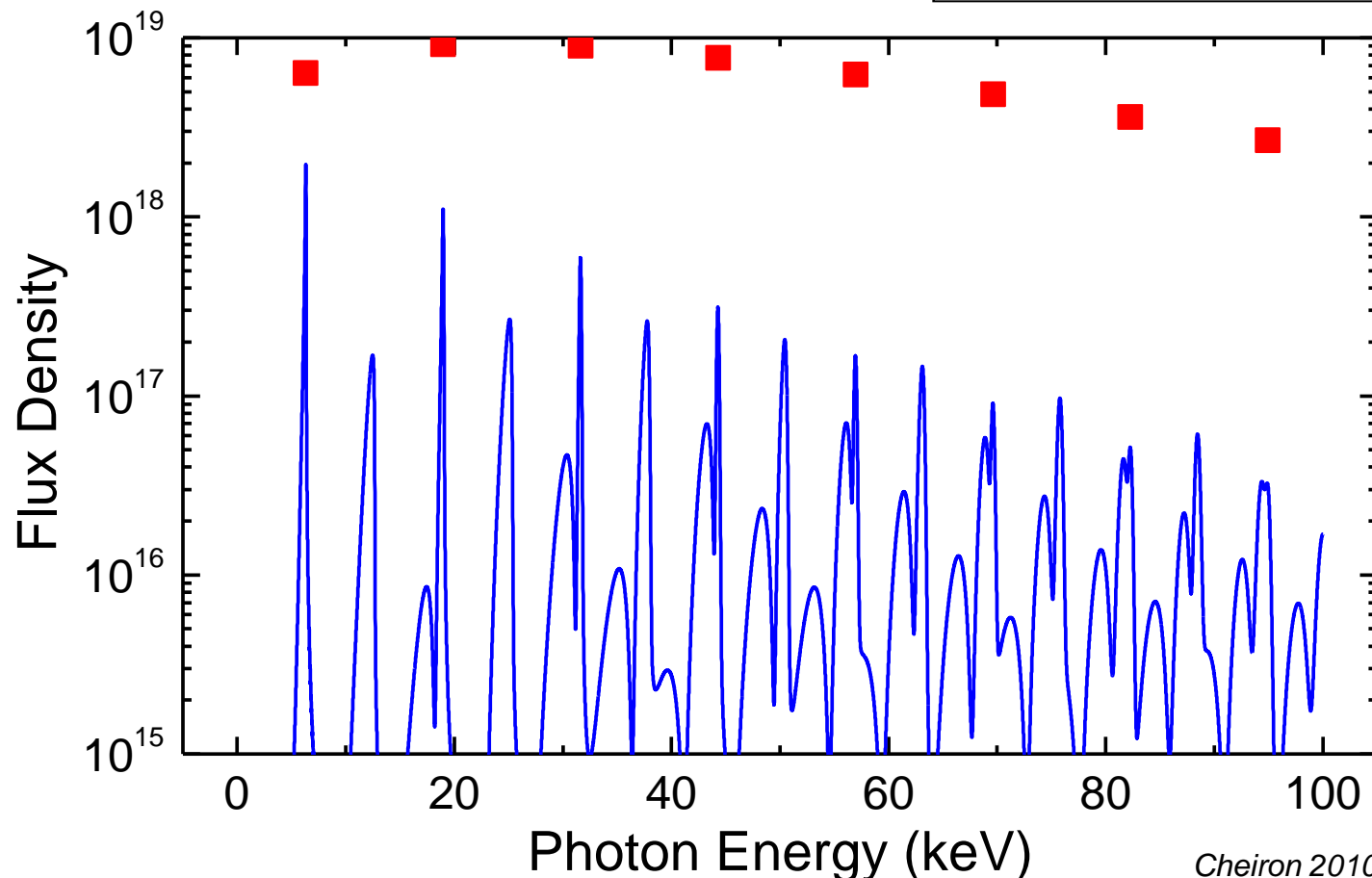
# Effects on the Higher Harmonics

Optical Emittance of UR:  $\lambda/4\pi$

Bandwidth of UR:  $\sim 1/nN$



Effects due to the  $e^-$  beam are larger for higher harmonics



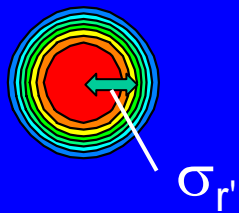
# Evaluation of Optical Properties

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- Optical properties of UR from the e-beam:
  - Calculate the UR from a single electron based on the theory of electrodynamics.
  - Integrate over the electrons in the beam (convolution)
    - Requires complicated numerical computation with a large number of parameters
    - Dedicated computer software (SPECTRA, SRW,...) is available
- Easy evaluation with Gauss approximation
  - Source size, angular divergence
  - Flux density, brilliance

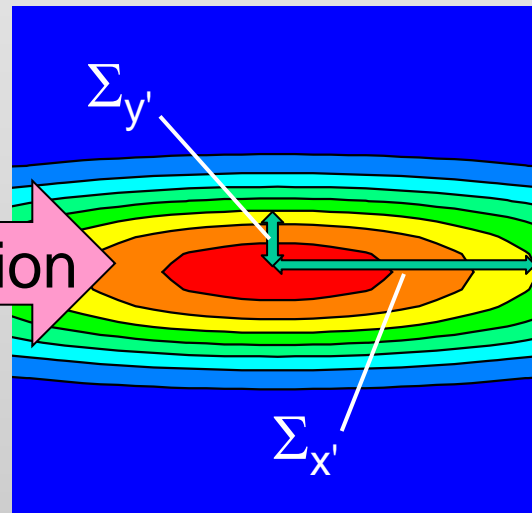
# Effective Size & Divergence

Angular profile of UR  
from a single electron

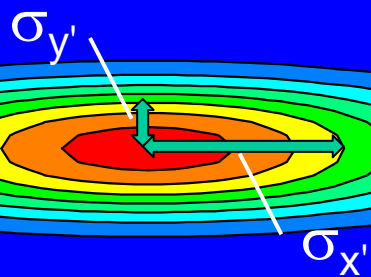


Convolution

$\Sigma_{y'}$



Angular profile of UR  
from the e-beam



Angular profile of the e-beam

By Gauss approx. and  
convolution theorem,

$$\Sigma_{x',y'} = \sqrt{\sigma_{r'}^2 + \sigma_{x',y'}^2}$$

Effective Angular Div.

$$\Sigma_{x,y} = \sqrt{\sigma_r^2 + \sigma_{x,y}^2}$$

Effective Source Size

# Effective Flux Density and Brilliance

$$\int_{-\infty}^{\infty} G \exp(-x^2/2\sigma^2) dx = G \times \sqrt{2\pi}\sigma$$

Effective width of a Gauss function is  $\sqrt{2\pi}\sigma$

**Total Flux**  $F = \boxed{\left. \frac{d^2 F}{dx' dy'} \right|_0} \times 2\pi\sigma_{r'}^2$   
 on-axis flux density with zero-emittance beam

**Effective Flux Density**  $\left. \frac{d^2 F}{dx' dy'} \right|_e = \frac{F}{2\pi\sigma_{x'}\sigma_{y'}} = \left. \frac{d^2 F}{dx' dy'} \right|_0 \frac{\sigma_{r'}^2}{\sigma_{x'}\sigma_{y'}}$

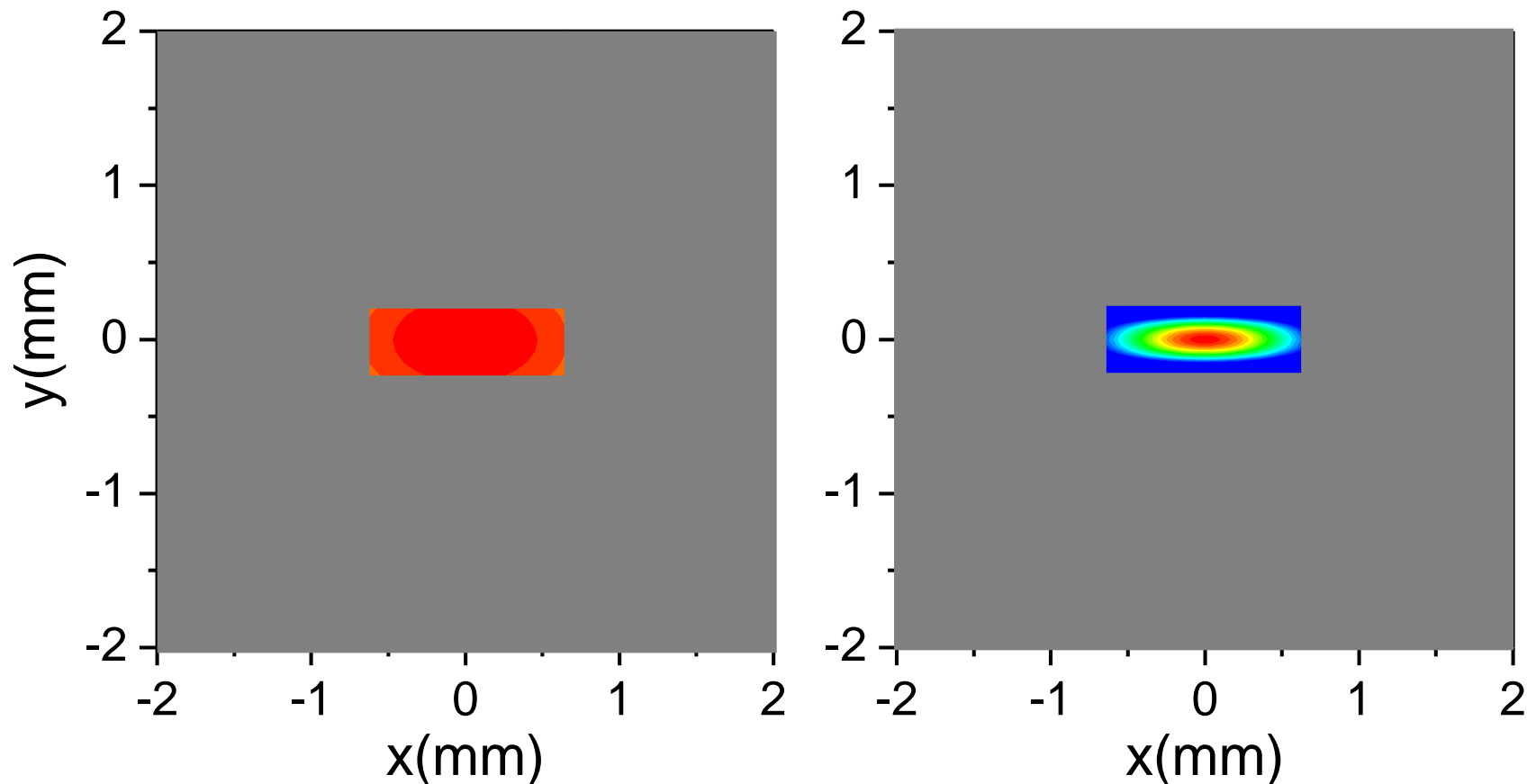
**Effective Brilliance**  $B_e = \frac{F}{4\pi^2\sigma_x\sigma_y\sigma_{x'}\sigma_{y'}} = \left. \frac{d^2 F}{dx' dy'} \right|_0 \frac{\sigma_{r'}^2}{2\pi\sigma_x\sigma_{x'}\sigma_y\sigma_{y'}}$

# Heat Load on Optical Elements

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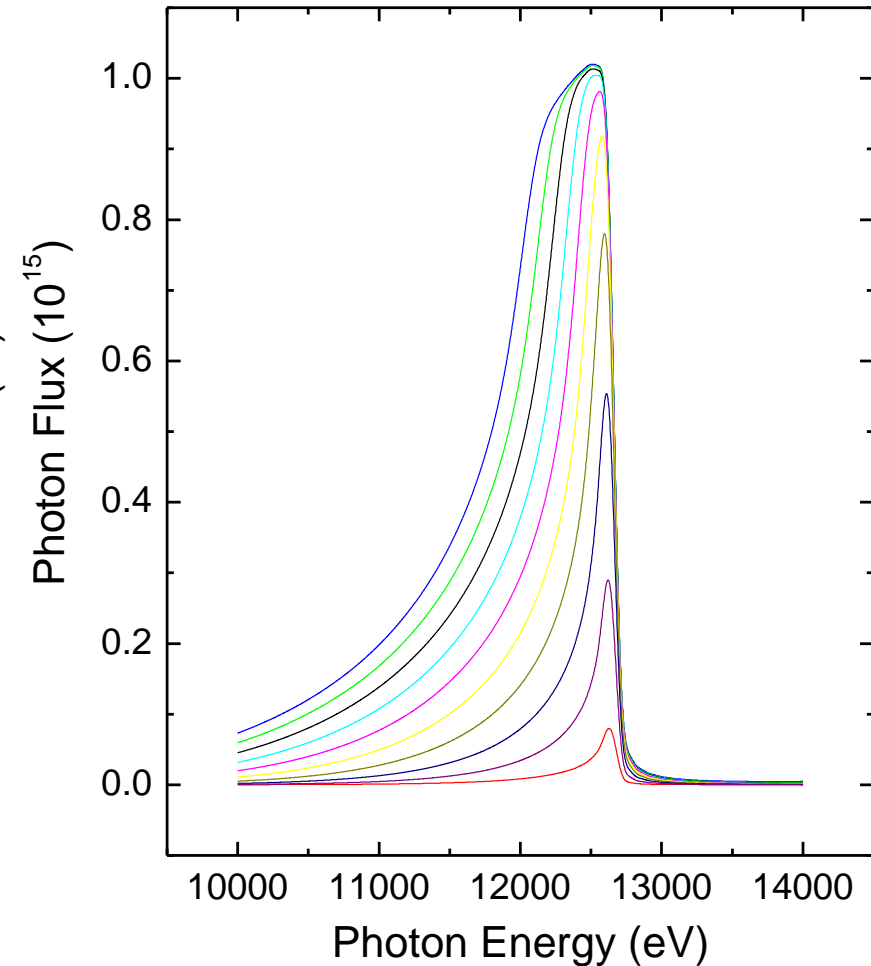
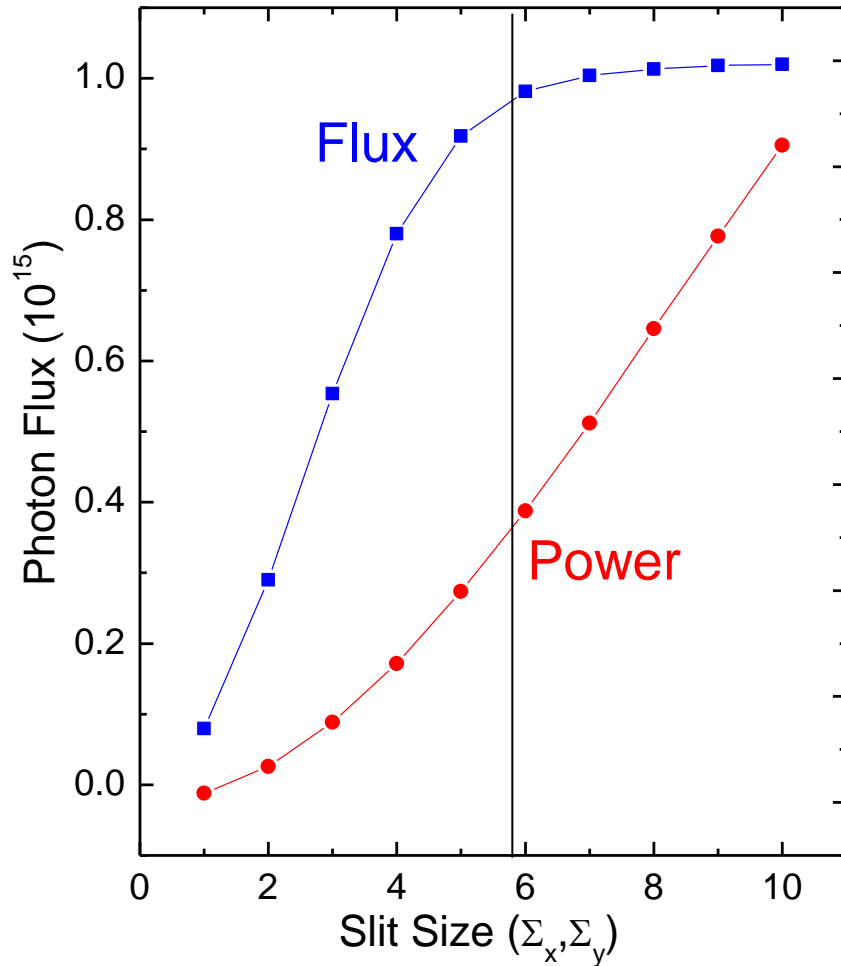
- SR emitted from the light source is processed by several optical elements before irradiation to the sample, such as the focusing mirror, monochromator.
- These elements can be easily damaged by the heat load brought by the SR.
- It is thus important to reduce the heat load as much as possible without sacrificing the flux by means of the XY slit at the front-end section.

# Spatial Profile of Power and Flux



The power profile is much broader than the flux. Extraction of SR with an appropriate slit significantly reduces the heat load.

# Optimum Slit Size?



Too wide aperture results in broader bandwidth and unwanted heat load.

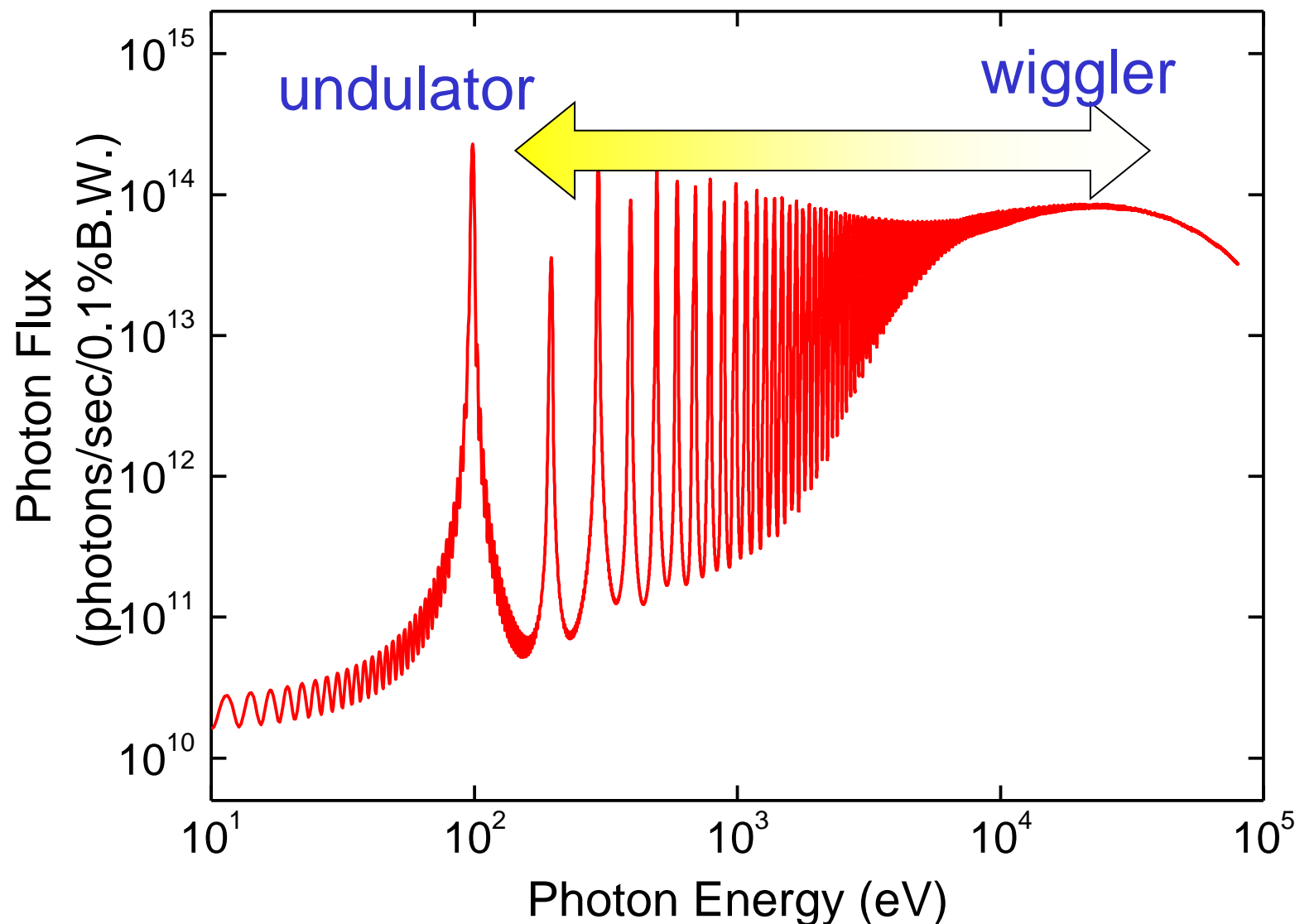
# Wiggler? Undulator? (1)

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- Wigglers are identical to undulator from the point of view of magnetic circuit.
- It is generally said that the  $K$  value distinguishes between the two, however, this is not exactly correct.
- What we should take care is the region of photon energy to be utilized for application.



# Wiggler? Undulator? (2)



# Undulator Radiation Gallery

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- For quantitative evaluation of SR, a computer code “SPECTRA”, which has been developed and maintained in SPring-8 is available.
- SPECTRA also offers a function to “visualize” the computation results for further understanding of SR.
  - brilliance curve & spectrum
  - on- and off-peak angular profiles of flux
  - on- and off-axis spectra
  - effects of opening the slit aperture
  - undulator-to-wiggler transition

# Other Topics Not Addressed

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- Quantitative descriptions of SR
- Light sources for circular polarization and schemes for fast helicity switching
  - helical undulator & elliptic wiggler
  - chicanes&choppers, kicker magnets
- Effects on the electron beam
  - natural focusing
  - beam-axis fluctuation due to COD variation
- R&Ds toward shorter magnetic period
  - superconducting undulators
  - cryogenic permanent magnet undulators
- Coherent SR for intense THz light
- Undulators for SASE-based X-ray FEL